Modification of molecular packing: crystal structures and magnetic properties of monomeric and dimeric difluorophenyl-1,2,3,5-dithiadiazolyl radicals

Arthur J. Banister,^a Andrei S. Batsanov,^a Owen G. Dawe,^a Penelope L. Herbertson,^a Judith A. K. Howard,^a Sean Lynn,^b Iain May,^a J. Nicholas B. Smith,^c Jeremy M. Rawson,^{*,c} Toni E. Rogers,^d Brian K. Tanner,^d Guillermo Antorrena^e and Fernando Palacio^e

^a Department of Chemistry, The University of Durham, South Road, Durham, UK DH1 3LE

^b Glaxo-Wellcome Research Centre, Gunnels Wood Road, Stevenage, Hertfordshire, UK SG1 2NY

^c Department of Chemistry, The University of Cambridge, Lensfield Road, Cambridge,

^d Department of Physics, The University of Durham, South Road, Durham, UK DH1 3LE

^e Instituto de Ciencia de Materiales de Aragon, CSIC-Universidad de Zaragoza, E-50009, Spain

Two dithiadiazolyl radicals of formula $[F_2C_6H_3CNSSN]^*$ were prepared and characterised by X-ray crystallography; in the solid state the 2',3'-difluoro isomer forms discrete twisted dimers through a spin-paired singly occupied–singly occupied molecular orbital interaction with an intra-dimer S · · · S distance of 3.020(4) Å, whereas the 2',5'-isomer is composed of uniform stacks with unexpectedly long intra-stack S · · · S contacts [3.544(3) Å].

The majority of dithiadiazolyl radicals, [RCNSSN], lose their paramagnetic nature in the solid state through an out-of-plane, spin-paired, dimerisation process in which the separation of the heterocyclic rings within the dimeric unit is typically 2.99-3.14 Å;¹ intermediate between twice the covalent radii² ($\hat{2}$.10 Å) and twice the sum of the van der Waals radii³ (4.06 Å). The energy of this interaction (*ca.* 35 kJ mol⁻¹)⁴ is much less than that for a covalent bond (*cf.* 265 kJ mol⁻¹ for S–S)⁵ but is extremely large on the scale of intermolecular interactions. Therefore counteracting this dimerisation process by means of weak van der Waals forces (the principle instrument of 'crystal engineering') is a difficult task. Other workers have shown that hydrogenbonding interactions can be successfully used to control molecular packing⁶ (e.g. in hydroxy-substituted nitronylnitroxide radicals⁷), but such an approach is impossible in these compounds because dithiadiazolyl radicals readily react with acidic protons.^{†,8} Recently we found the first dithiadiazolyl radical, $[p-NCC_6F_4CNSSN]$, which retains its monomeric nature in the solid state.^{9,10} This radical has been found in two morphologies, 9,10 both of which are paramagnetic. The $\beta\text{-phase}$ is exceptional because it undergoes a phase transition to a weakly ferromagnetic state at 36 K.¹⁰ We believe that the $F \cdots F$ repulsions in a head-to-head dimer diminish energy gain on dimerisation to such an extent that the favourable in-plane $CN \cdots S$ interactions in [NCC₆F₄CNSSN]' can overcome it.⁹ These results illustrate that we are able to overcome the spin-paired dimerisation process to produce radicals exhibiting long-range magnetic order, and that fluorinated aryl substituents appear to play a prominent role in determining the solid-state structure. Herein we report the structures of two novel difluorinated



dithiadiazolyl radicals; $[2',3'-F_2C_6H_3CNSSN]$ [•] I and $[2',5'-F_2-C_6H_3CNSSN]$ [•], II, each of which possesses an unusual solid-state architecture.

Compounds **I** and **II** were prepared from 2,3- or 2,5difluorobenzonitrile respectively, using standard synthetic procedures¹¹ and crystals suitable for X-ray diffraction studies \ddagger were grown by vacuum sublimation [10⁻² Torr (Torr \approx 133.322 Pa), 80–120 °C].

Molecules **I** and **II** both contain planar dithiadiazolyl rings of unexceptional geometry. The asymmetric unit of **I** comprises two molecules. Within each molecule the dithiadiazolyl and benzene rings are nearly coplanar with a torsion angle around the C(1)–C(7) bond of 6°. Unexpectedly, the two molecules are associated in the solid state as a twisted dimer (Fig. 1) with a mean intra-dimer separation of *ca.* 3.1 Å between nearly coplanar dithiadiazolyl rings (dihedral angle 3.4°) and with the closest intra-dimer contact S(2A) · · · S(2B) of 3.020(4) Å. Previously such twisted conformations have only been observed in

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 $[\]dagger$ Recently the crystal structures of some dithiadia zolyl–metal complexes have shown ring protonation. $^{8b.c}$

[‡] Crystal data: $C_7H_3F_2N_2S_2$ 1, M = 217.23, triclinic, space group $P\overline{I}$, 4 Crystal data: C₇Tr₃r₂P₂S₂ **1**, *M* = 217.25, themic, space group *P*1, *a* = 6.637(6), *b* = 8.768(11), *c* = 13.546(5) Å, *a* = 88.79(6), β = 86.37(5), *γ* = 80.40(9)°, *U* = 776(1) Å³, μ = 6.14 mm⁻¹, *Z* = 4, *D_c* = 1.86 g cm⁻³, *T* = 150(2) K. Crystal data: C₇H₃F₂N₂S₂ **2**, *M* = 217.23, tetragonal, space group *P*4₂/*n*, *a* = 21.28(1), *c* = 3.544(3) Å, *U* = 1604(3) Å³, μ = 5.94 mm⁻¹, *Z* = 8, *D_c* = 1.80 g cm⁻³, *T* = 150(2) K. Data for **I** and **II** were collected on a Rigaku AFC6S four-circle diffractometer with graphitemonochromated Cu-K α radiation ($\lambda = 1.541$ 84 Å) in an ω scan mode. For I, 2365 reflections were measured $(32 < 2\theta < 40^{\circ})$ of which 2365 data were unique. For II, 2457 reflections were measured $(30 < 2\theta < 38^\circ)$ of which 1438 were unique. Both structures were solved by direct methods and refined using full-matrix least-squares refinements on F^2 values of all data using SHELXTL software.¹² At convergence; for **I**, $R1 = 0.070 [F > 2\sigma(F)]$, wR2 = 0.197 and the goodness of fit = 1.11 for all data and 236 parameters; for **II**, R1 = 0.090 [$F > 2\sigma(F)$], wR2 = 0.281 and the goodness of fit = 1.09 for all data and 136 parameters. All non-H atoms were refined anisotropically and H atoms were treated as 'riding'. The aryl group in II was found to be disordered over two orientations differing by a 180° rotation about the C(1)-C(7) axis; the site occupancies were refined to approximately 2:1. Maximum and minimum residual electron densities were within ± 0.9 for I and ± 0.7 e Å⁻³ for II. CCDC reference number 186/584.



Fig. 1 The asymmetric unit of compound **I** and its translational (along *x* axis) equivalent, showing the shortest $S \cdots S$ contacts. Selected bond lengths (Å) and angles (°): S(1A)-S(2) 2.105(3), S(1B)-S(2B) 2.095(3), S(1A)-N(5A) 1.636(6), S(1B)-N(5B) 1.634(7), S(2A)-N(3A) 1.617(6), S(2B)-N(3B) 1.163(7), N(3A)-C(7A) 1.36(1), N(3B)-C(7B) 1.35(1), N(5A)-C(7A) 1.32(1), N(5B)-C(7B) 1.32(1); N(5A)-S(1A)-S(2A) 93.4(2), N(5B)-S(1B)-S(2B) 93.6(2), S(1A)-S(2A)-N(3A) 95.2(3), S(1B)-S(2B)-N(3B) 95.0(2), S(2A)-N(3A)-C(7A) 113.6(5), S(2B)-N(3B)-C(7B) 114.3(5), S(1A)-N(5A)-C(7A) 115.3(6), N(3A)-C(7A)-N(5A) 122.4(6), N(3B)-C(7B)-N(5B) 121.8(7)

derivatives with sterically demanding substituents (Me, CF₃ or NMe₂)¹³⁻¹⁵ whereas planar aromatic substituents normally favour eclipsed (cisoid) configurations.¹ The driving force behind the twist in **I** is probably due to the asymmetric positions of the fluorine atoms which induce a comparatively large molecular dipole across the aromatic substituent, coupled with the intermolecular electrostatic interactions between the main centres of positive (sulfur) and negative (fluorine) charges [S(2A) ··· F(3B) and S(2B) ··· F(3A) at 3.03 and 3.07 Å respectively]. These intermolecular interactions occur between adjacent stacks of molecules and are close to the dithiadiazolyl ring plane.

Neighbouring dimers, related by the lattice translation *a*, pack so as to produce an infinite stack of dithiadiazolyl rings running parallel to the crystallographic *x* axis (Fig. 1). The mean inter-dimer separation (*ca.* 3.5 Å) and inter-dimer S····S contact 3.628(4) Å are considerably larger than the intra-dimer distances. The Peierls distortion which gives rise to this set of alternating short and long contacts between molecules is a common feature of the dithiadiazolyl ring system¹⁶ and leads to a dimeric structure which is correspondingly diamagnetic.

In comparison, molecule II exhibits a unique molecular packing arrangement for dithiadiazolyls. The tetragonal lattice consists of stacks of uniformly spaced dithiadiazolyl radicals (Fig. 2). The radical separation in II [equal to the length of the crystallographic c axis, i.e. 3.544(3) Å] is considerably larger than that normally observed ¹ for dithiadiazolyl radicals ($d_{s...s}$ 2.99-3.14 Å) although still less than twice the van der Waals radius of sulfur (4.06 Å when perpendicular to the bonds).³ Structurally, the radicals remain comparatively isolated and the capacity for sulfur atoms to undergo secondary interactions is compensated through strong inter-stack interactions. Four of these stacks are located around a 4, screw-axis in a 'pin-wheel' fashion, with the heterocyclic rings inclined by 8° to the screw axis (Fig. 3). At the centre of the pin-wheel, there are interstack S · · · S interactions [$d_{s...s}$ 3.493(3) Å] linking each molecule to its two nearest neighbours (above and below its own ring plane) in each of the two adjacent stacks. Further from the



Fig. 2 Molecular packing of compound **II** viewed down the [1 1 0] line, showing the uniform packing arrangement. Selected bond lengths (Å) and angles (°): S(1)–S(2) 2.087(3), S(1)–N(5) 1.627(7), S(2)–N(3) 1.633(7), N(3)–C(7) 1.361(9), N(5)–C(7) 1.314(11); N(5)–S(1)–S(2) 93.7(3), S(1)–S(2)–N(3) 95.3(2), S(2)–N(3)–C(7) 112.9(6), S(1)–N(5)–C(7) 115.7(5), N(3)–C(7)–N(5) 122.3(7)



Fig. 3 'Pin-wheel' packing of compound II viewed down the 4_2 screw axis (*z* axis)

centre of the pin-wheel there is a set of $S(1) \cdots N(3)$ and $S(1) \cdots F(6)$ contacts at 2.95 and 3.17 Å respectively (*cf.* corresponding sums of van der Waals radii³ at 3.20–3.63 and 2.90–3.38 Å).

Compound **II** provides the first example of an intrinsic, undistorted dithiadiazolyl stack; all other dithiadiazolyls reported to date undergo a Peierls distortion to form discrete dimer pairs (as in compound **I**). Peierls distortions in onedimensional systems tend to be suppressed by increasing the dimensionality of the system through the presence of significant inter-stack interactions^{16,17} (as in compound **II**). Oakley and co-workers¹⁸ have shown that this process can also be suppressed through iodine doping in [HCNSSN]·0.18I, **III**.

A few striking similarities between radicals **II** and **III** are worthy of further comment. In both structures the period of identity of the dithiadiazolyl stack coincides with that of the crystal lattice in the direction of the stack [3.544(3) and 3.352(5) Å for **II** and **III** respectively], *i.e.* the unit cell dimension (*c*) coincides with the intra-stack $S \cdots S$ separations. In both structures the stacks of dithiadiazolyl radicals are arranged around a high-order screw-axis in a pin-wheel motif (four-fold for **II** and six-fold for **III**) which is clearly favourable for creating a threedimensional network of inter-stack interactions. The six-fold arrangement in III forms a large channel (internal diameter 7.40 Å, twice the distance from the inner sulfur to the 6_1 axis) at the centre of the pin-wheel which is filled with randomly disordered I₂ molecules.¹⁸ In the case of II, the corresponding cavity is substantially smaller (internal diameter 4.26 Å, twice the distance from the inner sulfur to the 4_2 axis) and so it may be able to accommodate a lithium atom.

With the abnormally large intra-stack separation of dithiadiazolyl rings (ca. 0.5 Å greater than that normally observed for dimeric dithiadiazolyl radicals, and even larger than that observed in III), we suspected that II might be paramagnetic in the solid state. However, variable-temperature (5-300 K) magnetic measurements on II indicated that it was diamagnetic (diamagnetic susceptibility, $\chi_d = -5.55 \times 10^{-9}~m^3$ kg⁻¹) with a small paramagnetic contribution at low temperature equivalent to less than 0.1% of molecules with an unpaired spin.§

These results illustrate how small changes in molecular structure can lead to large changes in solid-state structure. Compound I provides the first example of electronically (rather than sterically) directed control of the dimeric structure normally associated with dithiadiazolyl radicals. Compound II is the first example of an undoped dithiadiazolyl radical which forms a uniform stack which does not suffer a Peierls distortion. Despite the unexpectedly long $S \cdots S$ contacts, II remains diamagnetic in the solid state. Further investigations are needed to determine at what point the separation of dithiadiazolyl rings becomes sufficient to induce paramagnetic behaviour.

 $\$ Oakley and co-workers have recently reported 19,20 several examples of dithiazolyl radicals in which the molecules form uniformly stacked structures with separations of ca. 3.5 and 3.7 Å. When the separation is around 3.5 Å the compound is essentially diamagnetic $^{\mbox{\tiny 19}}$ at room temperature (fraction of Curie spins at 300 K \approx 0.01) whereas with the larger intermolecular separation,²⁰ the fraction of Curie spins at room temperature is ≈ 0.3 . The essential diamagnetism in the former case is associated with a band gap of *ca.* 0.4 eV (eV \approx 1.602 18 \times 10⁻¹⁹ J) which gives rise to a filled valence band.

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